

## Characterization of a Spherical Proportional Counter in argon-based mixtures

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The Spherical Proportional Counter is a novel type of radiation detector, with a low energy threshold (typically below 100 eV) and good energy resolution. This detector is being developed by the network NEWS, which includes several applications. We can name between many others Dark Matter searches, low level radon and neutron counting or low energy neutrino detection from supernovas or nuclear reactors via neutrino-nucleus elastic scattering. In this context, this works will present the characterization of a spherical detector of 1 meter diameter using two argon-based mixtures (with methane and isobutane) and for gas pressures between 50 and 1250 mbar. In each case, the energy resolution shows its best value in a wide range of gains, limited by the ballistic effect at low gains and by ion-backflow at high gains. Moreover, the best energy resolution shows a degradation with pressure. These effects will be discussed in terms of gas avalanche properties. Finally, the effect of an electrical field corrector in the homogeneity of the gain and the energy threshold measured in our setup will be also discussed.

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## 1. Introduction

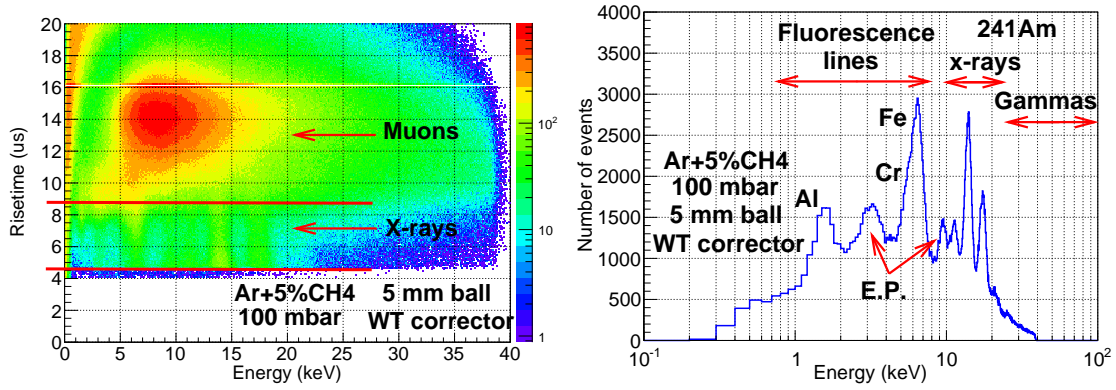
An active field in instrumentation for Particle Physics is the development of massive detectors with low background levels and low energy threshold. The range of applications is diverse: the search of dark matter in the form of Weakly Interacting Massive Particles (WIMPs) [1] or axions [2], the detection of low-energy neutrinos, radon and neutron counting. In this context, the Spherical Proportional Counter (SPC) was proposed in [3] as a candidate fulfilling all the requirements: a single readout channel reading a big gas volume (and mass); a potentially low intrinsic background level with a good selection of radiopure materials; and a high signal-to-noise ratio due to its low capacitance, which is proportional to the radius of the central electrode and does not depend on the vessel size. Indeed, previous works have shown that an energy threshold as low as 100 eV [4, 5] is feasible, keeping a good energy resolution (10% FWHM at 22.1 keV).

Even if there is only one channel, the pulse risetime can be used to discriminate complicated topologies like muons from point-like events like x-rays; or to set a fiducial volume, as the risetime depends on the radial position where the energy was deposited by gas diffusion. As an example of this feature, we present in figure 1 (left) the dependence of the risetime with the pulse amplitude for our SPC irradiated by a  $^{241}\text{Am}$  source covered by an aluminum foil. As the mean free path of photons is longer than the sphere's radius at low pressures, an x-ray may deposit its energy at any distance. As a result, it will create a distribution in risetime for a fixed amplitude in the 2D plot. This fact happens for the x-rays of the source (at 11.9, 13.9, 17.7 and 20.8 keV). In contrast, the fluorescence lines emitted from the stainless steel vessel (chromium at 5.5 keV and iron at 6.4 keV) and the foil (aluminum at 1.5 keV) form three spots at long risetimes. If we select the events whose risetime ranges in 4.5-8.7  $\mu\text{s}$  (red lines), we obtain the energy spectrum of figure 1 (right).

The advantages of SPCs have motivated several feasibility studies [5–9] and the groups interested in this technology have created the network NEWS (New Experiments With Spheres), which is now formed by institutes of France, Greece, China and Spain. In this context, the scalability of high gains and good energy resolutions to higher masses (pressures) is an open question, which is a key-point to increase the sensitivity of the detector to any particular application. This work tries to give a first answer, presenting the characterization of a SPC in two argon-based mixtures (with methane and isobutane) and for gas pressures up to 1250 mbar. The setup, the data-taking procedure and the analysis are described in detail in section 2. We also present the two configurations of the central rod (without and with a field corrector), whose main results will be respectively detailed in sections 3 and 4. Finally, we finish in section 5 with some conclusions and an outlook.

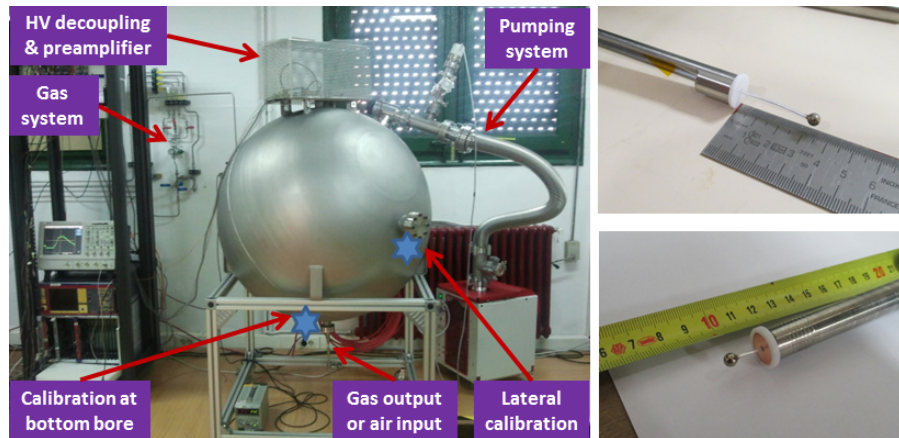
## 2. Detector description

The detector consists of a large spherical stainless steel vessel 1.0 meter in diameter and a small metallic ball 5 mm in diameter kept at the center by a stainless steel rod. The radiation from the calibration source or other sources (like muons and gammas) ionize the gas. As the central ball is set to a positive voltage and the vessel to ground, electrons drift to the ball where an intense electric field amplifies the charge. The charge movement induces a signal at the electrode, which is extracted by a teflon-based high voltage feed-through from the vessel. The signal is decoupled



**Figure 1:** Left: Dependence of the risetime with the pulse amplitude in Ar+5% methane at 100 mbar, obtained irradiating our SPC with a  $^{241}\text{Am}$  source whose 5.5 MeV alphas have been blocked by an aluminum foil. Muons show an event distribution centered at 8 keV and 14  $\mu\text{s}$ . Some muons events (see band below 3 keV at long risetime values) are not well analyzed as their risetime is longer than the analysis range. The x-rays from the source (at 11.9, 13.9, 17.7 and 20.8 keV) create different bands in risetime while the fluorescence lines emitted from the vessel (chromium at 5.5 keV and iron at 6.4 keV) and the aluminum foil (at 1.5 keV) form three spots at long risetime. The red lines delimit the fiducial volume for point-like events. Right: Energy spectrum after having selected the events in the fiducial volume. The Argon escape peaks (E.P.) of both fluorescence lines and  $^{241}\text{Am}$  x-rays lines are also present in the spectrum.

from the high voltage (powered by a CAEN N1471HA module) by a capacitance and then fed into a CANBERRA preamplifier, which is acquired by a Tektronix DS5054B oscilloscope. A view of the setup is shown in figure 2 (left), while on right caption, we show the two configurations used in these tests: in the first one the ball is situated at 30 mm of the rod, covered by a teflon cap; in the second one, the ball is 20 mm away from a copper plate, which works as a field corrector (generally called “umbrella”) and is isolated from the rod by a teflon cassette.

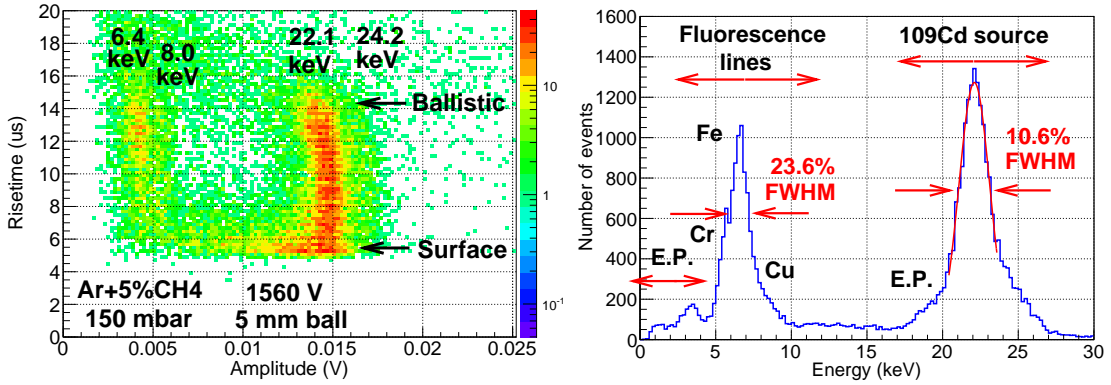


**Figure 2:** Left: A general view of the spherical proportional counter characterized in this work and the auxiliar facilities: the gas system, the pumping system and the in/output valves. The blue stars indicate the points where the cadmium source was installed, as described in the text. Right: the two configurations tested in this setup: a central ball without (top) and with (bottom) an umbrella field corrector.

The detector operates in seal mode, i.e., the vessel is first pumped out and then filled with an appropriate gas at a pressure from few tens of mbar up to 2 bar. Two mixtures have been studied: Ar+5% methane for pressures between 50 and 1250 mbar and Ar+2% isobutane for pressures between 100 and 400 mbar. No gas filter has been used in this work. The pumping lasted at least one night to reduce the outgassing rate from inner vessel's components below  $10^{-6}$  mbar l/sec. This value allows the operation of the detector during a week without any visible degradation.

The results presented here were obtained irradiating the SPC by a  $^{109}\text{Cd}$  source respectively situated inside the vessel at the bottom part for the rod without the field corrector (figure 2), and at the lateral bore for the second configuration. In the first case, we observed that the energy resolution was poor ( $\approx 30\%$  FWHM at 22.1 keV) when calibrating from the lateral bore. In contrast, values were good ( $\approx 10\%$  FWHM) when the SPC was calibrated from the bottom part. We concluded that the field was not homogeneous and that a field corrector was needed, as suggested in [3]. However, as described in section 4, the field became homogeneous only when the umbrella was set to negative values. This fact disagrees with the conclusions of the former publication. A simulation of the electric field is foreseen to evaluate the effect of the field corrector in our setup.

In a first step, signals acquired by the oscilloscope are smoothed. The first derivate is then calculated to define a temporal range where the Pulse Shape Analysis (PSA) is applied. This range is useful to avoid errors produced by pile-up events, unstable baselines or noise. Finally, the PSA calculates some pulse features like the baseline, risetime, or amplitude. As shown in figure 1, the risetime allows to discriminate x-rays from muons setting an upper limit.



**Figure 3:** Left: Dependence of the risetime with the pulse amplitude for a low voltage case in Ar+5% methane at 150 mbar, obtained irradiating the SPC with a  $^{109}\text{Cd}$  source. Right: A typical energy spectrum obtained after the risetime selection explained in text.

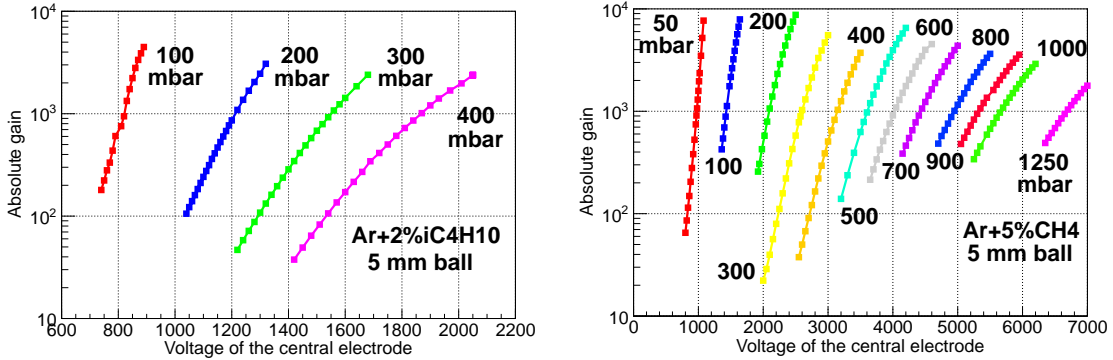
Two other effects have been observed: one is a decrease in amplitude for events with long risetimes at low gains (figure 3, left). This effect is caused by the ballistic effect, i.e., the time needed to collect all charge is comparable to the discharging time of the preamplifier. To compensate this degradation, an upper limit in risetime of  $13 \mu\text{s}$  has been applied. The second one is a population of 22.1 keV events with lower amplitudes at short risetimes (same figure as before). This effect happens only at the configuration without any field corrector. We have attributed these events to

x-rays depositing its energy near the ball, where a surface defect or a field distortion may reduce the detector's gain. To remove this contribution, a lower limit at risetime of 5  $\mu$ s has been set.

The resulting spectrum after the risetime selection is shown in figure 3 (right). The  $K_\alpha$  (22.1 keV) and  $K_\beta$  lines (24.9 keV) from Ag fluorescence are clearly distinguished, as well as the fluorescence lines of the vessel (crome at 5.5 keV, iron at 6.4 keV) and the source container (copper at 8.0 keV), and the corresponding Argon escape peaks (E.P.). The peak at 22.1 keV is used to obtain the gas gain and energy resolution, fitting it to a gaussian function (red line). The energy resolution of 6.4 keV line is around 23% FWHM for all optimum cases and configurations, near the expected value (20% FWHM) derived from the 22.1 keV peak. The difference between both values could be explained by fitting errors due to the presence of other peaks in the same range of energy.

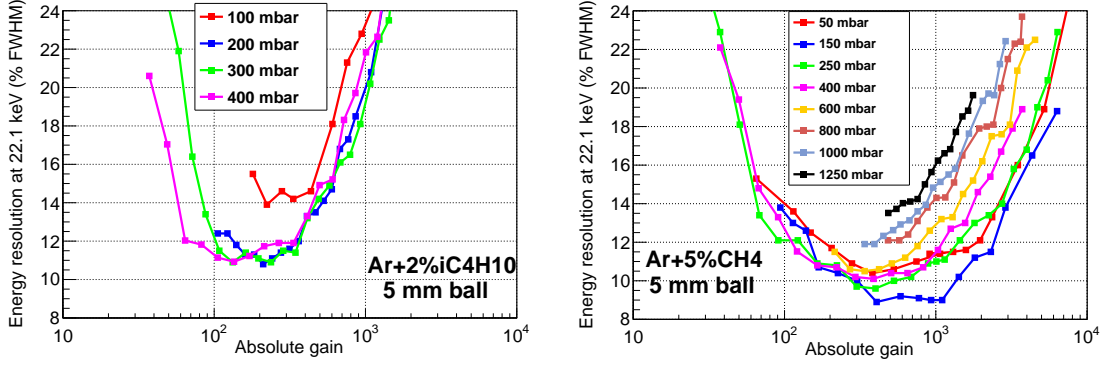
### 3. Results without the field corrector

The dependence of the peak position with the voltage generates the gain curves, shown in figure 4 for both argon-based mixtures. The gain reaches values higher than  $2 \times 10^3$  in all cases. Moreover, the spark limit was not observed in any case and the data-taking was resumed when the energy resolution degraded. The dependence of the gain with the voltage follows the expected exponential relation based on the Rose-Korff model [10], except for high voltages where gain tends to saturation. We have attributed this effect to ion-backflow [11], i.e, ions produced in the avalanche enter in the drift volume and produce distortions due to space-charge effects.



**Figure 4:** Dependence of the absolute gain with the voltage of the central electrode for Ar+2% isobutane (left) and Ar+5% methane (right) and gas pressures between 50 and 1250 mbar.

The energy resolution depends on the detector's gain, as shown in figure 5. At low values, the energy resolution degrades because the signal is comparable to noise. As the gain increases, it reaches its best value and stays constant during a range of gains, worsening again at higher gains. The surprising fact is how different are the gains where this last degradation starts:  $\approx 3 \times 10^2$  for Ar+2%iC4H10 and  $\approx 10^3$  for Ar+5%iCH4, i.e., a factor 3 better for methane even if isobutane quenches better the photons generated in the avalanche [12]. For this reason, we can not attribute this effect to an increase of gain fluctuations or the proximity to the spark limit, but to an effect probably correlated with the former gain saturation. Just as a comparison, energy resolution starts degrading in argon-based mixtures at  $10^4$  for Micromegas detectors [13].



**Figure 5:** Dependence of the energy resolution (% FWHM) at 22.1 keV with the absolute gain for Ar+2% isobutane (left) and Ar+5% methane (right) and gas pressures between 50 and 1250 mbar.

The best values obtained for energy resolution at 22.1 keV are 9% FWHM in Ar+5% methane at 150 mbar and 11% FWHM in Ar+2% isobutane at 200-400 mbar, as shown in figure 5. As noted before, better values were expected for isobutane as this gas is a better quencher. These values are similar to those reported in [4]. Finally, we have observed a degradation of energy resolution with pressure for the argon-methane mixture. Its physical origin is still unknown but similar tendencies have been reported for Micromegas detectors in  $\text{CF}_4$  [14] and Xe-TMA [15].

#### 4. First results of the field corrector

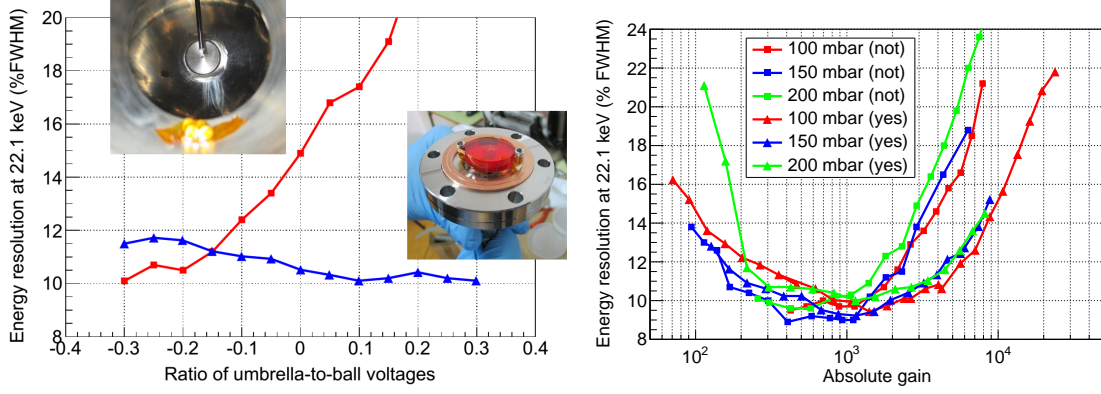
The effect of the umbrella field corrector (figure 2, top-right) was studied calibrating the SPC alternatively from the lateral bore and the bottom bore (blue stars at figure 2, left) in Ar+5% methane at pressures between 100 and 200 mbar. The gain increased if negative voltages were applied to the umbrella for both cases but no maximum was found. In terms of energy resolution (figure 6, left), we observed a clear improvement for lateral calibrations if the field corrector was set to negative values, reaching an optimum at a ratio of -0.2. However, the energy resolution slightly degraded for bottom calibrations. As a compromise, we fixed the ratio of umbrella-to-ball voltages to -0.2 when generating gain curves. Then, we replaced the lateral cap by an aluminized mylar window of  $3.5 \mu\text{m}$  thickness, and verified that the gain was uniform in all radial directions, because the vessel's fluorescence lines of a long run (see figure 1) were clearly defined.

The absolute gain was a factor 5 higher in this configuration, i.e., a lower voltage must be applied to reach the same gain. Apart from that, the degradation observed in energy resolution at high gains appeared at higher values (figure 6, right), which allowed us to reach an energy threshold below 1 keV (figure 1, right). This value is still far from the 0.1 keV reported in [4] because the DAQ used in this work has no trigger optimization. We expect to implement this type of electronics in future updates. The best value for energy resolution was 9.5% FWHM at 22.1 keV for 100 and 150 mbar, similar to the configuration without the field corrector.

#### 5. Conclusions and prospects

The Spherical Proportional Counter (SPC) is a novel type of radiation detector, with a low





**Figure 6:** Left: Dependence of the energy resolution (% FWHM) at 22.1 keV with the ratio of umbrella-to-ball voltages for the electrode configuration with the field corrector calibrating the SPC by a <sup>109</sup>Cd source alternatively situated inside the vessel at the lateral bore (red squares) and the bottom part (blue triangles). Data has been taken in Ar+5% methane at 100 mbar and setting the central ball at 1350 V. The two figures at the plot show how the <sup>109</sup>Cd source was installed. Right: Dependence of the energy resolution (% FWHM) at 22.1 keV with the absolute gain in Ar+5% methane and gas pressures between 100 and 200 mbar for the configuration without (squares) and with the field corrector (triangles).

energy threshold and good energy resolution. This detector has a wide range of applications like Dark Matter searches, neutrino or neutron detection. To increase the sensitivity of this detector, the scalability of its good properties is crucial. For this purpose, we present the characterization of a 1 meter diameter SPC in two argon-based mixtures and for pressures up to 1250 mbar. We have studied two rod configurations: without and with an umbrella field corrector.

For the first configuration, gains as high as  $2 \times 10^3$  have been reached. Gain curves show a saturation effect at high values, which seems to be correlated with a degradation in energy resolution and has been attributed to ion-backflow. The best values obtained for the energy resolution at 22.1 keV are 9% FWHM in Ar+5% methane at 150 mbar and 11% FWHM in Ar+2% isobutane at 200-400 mbar respectively for gains  $(0.2 - 1.0) \times 10^3$  and  $(0.1 - 0.2) \times 10^3$ . The energy resolution also degrades at high pressure but its physical origin is still unknown. In the second case, the umbrella field corrector has made uniform the detector's gain in all radial directions. The best value for energy resolution is 9.5% FWHM at 100-150 mbar. Energy resolution also degrades at high gains, but at higher values than in the first case, which has allowed us to reduce the energy threshold below 1 keV. The new electrode is now being characterized for pressures up to 1250 mbar.

In near term, we plan to update the acquisition system to further reduce the energy threshold and to study other light gases like neon and helium. Our studies with these gases can be applied in the project SEDINE, developed by part of the NEWS network. This project is operating a radiopure copper spherical vessel at the Modane Underground Laboratory (LSM) since end 2012, and aims at building a sphere of 2 meters of diameter filled with a light gas at 10-20 bar for Dark Matter searches. This future experiment may reach sensitivities near DAMA signal for WIMPs masses as low as 1 GeV if background level is around  $10^{-2} \text{ keV}^{-1} \text{ kg}^{-1} \text{ day}^{-1}$  and the energy threshold is kept below 0.1 keV for high pressures. For further details, the reader is referred to [6, 5].

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